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GAMMA-RAY LASER CANDIDATES

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Extension of nuclear structure data base searches for gamma-ray laser candidates

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Results from a data base search of computerized nuclear structure libraries have been extended and augmented so as to expand the information available for nuclei suitable as gamma-ray laser candidates. The spectrum of nuclear levels occurring in deformed rotational nuclei have been calculated and have been used in conjunction with isomeric state data for odd-A systems. The results of this augmentation effort are presented with particular emphasis on results obtained for ^{177}Lu , ^{179}Hf , and ^{235}U . For these cases some possibly interesting cases were identified that meet energy spacing criteria. However, significant advance factors exist for them which negate their interest for gamma-ray laser applications.

1. INTRODUCTION

In a previous computerized search¹ of experimentally based nuclear structure libraries,^{2,3} we identified eight pairs of nuclear levels that could be appropriate for gamma-ray laser applications. Specifically, such pairs consist of a long-lived nuclear isomer which acts as a laser storage state) lying close (within several hundred electron volts) to a short-lived state. If two such states exist, and the angular momentum difference between them is not overly great, then an external radiation source could be used to effect a transition between them, thus producing depopulation that could eventually lead to lasing.

In addition to identification of these level pairs, the search of Ref 1 produced a tabulation of long-lived isomeric states for odd-A and odd-odd nuclei. For such identified cases, some possibility exists that short-lived levels could exist nearby that heretofore had not been measured experimentally. Thus, if one were to augment this present experimental level data with results from theoretical calculations, then the expanded set comprised of experimentally and theoretically derived levels may provide new cases in which interesting gamma-ray laser candidates could be found.

We report in this paper the beginning of the process of level data augmentation using theoretical nuclear structure model results. We have concentrated on permanently deformed nuclei (rare earth and actinide) where simple models applicable to odd-A systems are appropriate. In the following sections the model will be described, nuclei appearing in the

isomeric state list of Ref.1 for which this model is appropriate will be identified, and results will be presented for closely spaced level pairs determined in this study.

2. APPROACH

The methods employed utilize the particle-rotor model of Bohr and Mottelson,⁴ which is applicable to deformed nuclei and which exploits symmetry properties arising from the deformations. The nuclear shape degrees of freedom of such nuclei lead to rotational bands of nuclear levels which are built upon the intrinsic states of the system. Thus, if one can specify the parameters of each rotational band, then one can construct the energy spacings, spins, and parities of the individual levels occurring in the band. We accomplish this using the particle-rotor model. The procedure (to be described in more detail later) allows one to calculate, given a relatively small amount of data, additional nuclear levels which may be missing in existing experimental nuclear structure data bases. Such information can then be combined with existing level data to identify cases that are possibly interesting for gamma-ray laser applications.

In our application of this technique, the first step was the identification of nuclei associated with the 51 odd-A isomeric states, listed in Ref 1, for which this model would be applicable. Of the 51 nuclei presented there, eleven can be described using a simple version of the particle-rotor model. These are the rare earth nuclei, ^{153}Ho , ^{165}Dy , ^{167}Lu , ^{177}Lu , ^{177}Hf , ^{179}Hf , ^{179}W , ^{187}Os , ^{187}W , and the actinide nucleus, ^{235}U . After identification of a nucleus as appropriate for the particle rotor model, a second criterion must be met, namely the isomeric state of interest must lie at a high enough excitation energy that a real possibility exists for new unmeasured levels to be identified. In our application these must lie close to the ground state. An extreme example, ^{235}U , appears in Fig.1 that illustrates this point. Here the isomeric state of interest (which is the first member of the excited $1-2[611]$ band) lies only 100 electron volts above the ground state band. The spacing of the higher lying levels in this band (the $9/2^-$, $11/2^-$, $13/2^-$ states, etc.) is on the order of tens of kiloelectron volts. Any higher-lying band members identified through theoretical calculations would lie at even higher excitation energies. Thus, the isomeric levels (other than the ground state) lie within an energy spacing that would be of interest for a gamma-ray laser application.

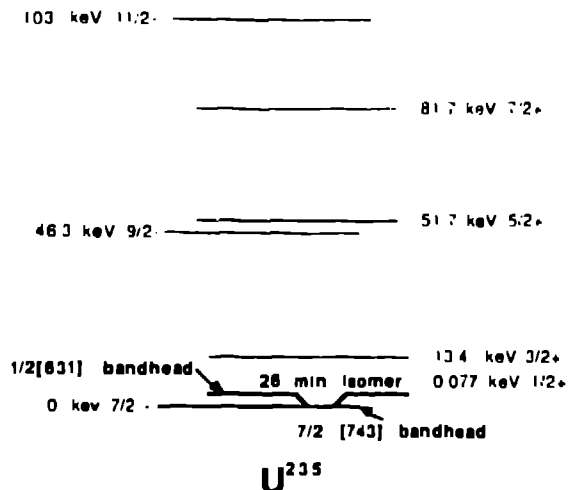


Figure 1. The band and level structure of ²³⁵U. Energies of the 7/2[743] ground-state band members appear on the left; energies of the 2(631) band, which includes the isomer, appear on the right.

After consideration of this second criterion only 3 of the 11 nuclei have their known isomeric levels occurring at high enough excitation energy for there to be a possibility of missing levels can be identified. These are ¹⁷⁷Lu ($E_x = 970.2$ keV, $J^\pi = 23/2^-$, half-life = 160.5 days), ¹⁷⁷Hf ($E_x = 2740$ KeV, $J^\pi = 7/2^-$, half-life = 51.4 min) and ¹⁷⁹Hf ($E_x = 05.7$ keV, $J^\pi = 25/2^-$, half-life = 25.1 days). For these nuclei, rotational bands were identified from existing experimental data and were constructed using the following particle-rotor expression to complete the spectrum of clear levels for each band:

$$E(J,K) = \epsilon_K + (\pi^2/2I) [J(J+1) - K^2 + \delta_K \frac{1}{2}(-1)^{J+1/2} a(J+1/2)], \quad (1)$$

where ϵ_K is the single particle (quasi-particle) energy, I is the moment of inertia, J is a total angular momentum, K its projection on the body-fixed symmetry axis, and a is the coupling parameter. The last term in Eq. (1), in which a appears, arises from Coriolis interactions. It has been assumed that the so-called recoil term can be absorbed into the single-particle energy. Finally, band mixing terms have been ignored.

The solution of Eq. (1) then requires knowledge of three states to determine the three unknowns (ϵ_K , I , a) for $K=1/2$ and knowledge of two states to determine the two unknowns (ϵ_K , I) for $K=1/2$. If insufficient experimental information exists (the usual case), one can use Nilsson⁶ or Nilsson-like calculations to complete, as much as possible, a low-lying spectrum. Note that the first members of a given low-lying band are the

most accurately known experimentally. In both cases the best possible values are obtained for ϵ_K , I , and a (if $K=1/2$) provided that one does not attempt to calculate band members too high in spin values where $I=I(\omega)$. Another reason to confine the calculation of the discrete level spectrum to lower excitation energies is that more complex states than those treated here arise as the excitation energy increases. The code NUKLEV was developed to carry out solutions of Eq. (1) as just described.

Table 1 lists for the three cases considered here the rotational band information (bandhead energies, spin, and K quantum number) used to produce the results described in the next section. This information was taken from data appearing in Ref 6-7.

Table 1. Band Parameters

	Band Energy (keV)	K	Bandhead State Spin-Parity
¹⁷⁷ Lu	0	7/2	7/2 ⁺
	150.39	9/2	9/2 ⁺
	457.9	5/2	5/2 ⁺
	569.62	1/2	1/2 ⁺
	761.65	1/2	1/2 ⁺
	970.15	23/2	23/2 ⁻
	1230.7	11/2	11/2 ⁺
	1356.9	15/2	15/2 ⁺
	1502.6	13/2	13/2 ⁻
¹⁷⁷ Hf	0.	7/2	7/2 ⁺
	321.3	9/2	9/2 ⁺
	508.1	5/2	5/2 ⁺
	559.4	1/2	1/2 ⁺
	608.	1/2	1/2 ⁺
	745.9	7/2	7/2 ⁺
	805.7	1/2	1/2 ⁺
	1057.8	1/2	1/2 ⁺
	1315.4	27/2	27/2 ⁻
	1434.	3/2	3/2 ⁺
1634.	1/2	1/2 ⁺	
1882.	1/2	1/2 ⁺	
¹⁷⁹ Hf	0.	1/2	1/2 ⁺
	214.	7/2	7/2 ⁺
	374.8	1/2	1/2 ⁺
	518.4	5/2	5/2 ⁺
	614.	1/2	1/2 ⁺
	720.	3/2	3/2 ⁺
	870.2	1/2	1/2 ⁺
	1305.7	25/2	25/2 ⁻

The next step in this process will involve similar procedure applied to a subset of the isomeric states identified for odd-odd nuclei in Ref 1. Here the chances of successful identification of potentially interesting level pairs may be increased significantly due to the complexity of low-lying band structures in such nuclei. In these instances, bandhead identification may occur through the examination of systematic trends obtained by studying the coupling of single-particle orbits in adjacent odd-A nuclei or through use of Nilsson-like particle orbits. In isolated cases identified so far, theoretical evidence has been presented for the existence of levels that are close to low-lying isomers, as was done in the case of ^{158}Ho by Sood et al.⁸ These investigations will also require information obtained from more microscopic nuclear structure development currently underway⁹ that allows more realistic specification of contributions due to residual n-p interactions.

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3. RESULTS

In this section results obtained for ^{177}Lu , ^{177}Hf , and ^{179}Hf are presented. For ^{177}Lu , nine rotational bands were used that resulted in a total of 51 calculated levels extending to an excitation energy of 1750 keV. Of these, the closest level to the 160-ray isomer at $E_x = 970.15$ keV was the $5/2^-$ level lying at 962 keV, which has been identified previously experimentally. No new additional theoretical levels were identified at energies below that of the isomeric state.

For our investigation of ^{177}Hf , data for twelve bands were used in the calculation, a condition necessitated in part by the high excitation energy (2740 keV) of the isomer. This calculation produced 127 levels up to an excitation energy of 2787 keV as compared with approximately fifty levels that are known experimentally. The closest calculated levels to the $37/2^-$ isomer are a $25/2^-$ level at 2709 keV and a second $25/2^-$ level at 2743 keV. Thus, the closest level is predicted to lie several kilovolts from the isomer. Even more important are the large changes in spin and K quantum number required to induce a transition between the isomer and the closest short-lived level. In this case, $\Delta J = 6$ (an E6 transition) and $\Delta K = 18$ exist, which produce overwhelming hindrance factors for transitions between these levels.

Finally, for ^{179}Hf , eight rotational bands were used to produce 47 levels up to an energy of 1366 keV vs the 26 that are known experimentally. The closest levels to the $25/2^-$ isomer at 1105.7 keV are the $17/2^-$ state calculated to occur at $E_x = 1105.3$ keV and the $11/2^-$ state at 1131 keV. Although the energy difference between the isomer and the state calculated to exist at 1105.3 keV is 400 electron volts (which is attractive from the point of view of the gamma-ray laser criteria discussed earlier), once again large spin and K differences will effectively eliminate transitions between these two level pairs. In this case, an E4 transition and a change in nine units of K would be required to effect a transition.

4. CONCLUSIONS

A simple particle-rotor model has been applied to augment experimental level data used in searches for nuclear level pairs appropriate to gamma-ray laser concepts. Two significant problems hampered this effort. The first occurred because the experimentally determined odd-A isomeric levels of interest lay, for the most part, at low excitation energies below the energy region where additional levels would be predicted. A more fundamental problem exists in that, even when the level spacing criteria of less than a kilovolt condition was satisfied, unacceptably large differences in spin and K quantum number required between the members of the pair. These in turn produce large hindrance factors for transitions between the two levels.